Thio and Imino Chelates

we obtained $\beta_1 = 1.0 \pm 0.5$ for the formation constant of $Eu(SCN)²⁺$ in agreement with the results of solvent extraction³ and UV spectroscopy. In the complex, the C-S stretching band is shifted to higher wavenumbers compared to free SCN-. This behavior has been correlated with the coordination of SCN- through the nitrogen.28 When **S** is the donor atom, the C-S stretching is observed at lower wavenumbers.

IV. Conclusion

The NMR, UV, and Raman experiments reported in this work concur to show that the binding of thiocyanate ions to trivalent lanthanide ions gives rise to inner-sphere complexes. The NMR spectroscopy indicates in particular that the spin density delocalized from the metal ions to the directly bonded nitrogen atom of the SCN- ligand is of the same order as that found for oxygen in the water solvation molecule^.'^ **A** similar situation was found in the case of cobalt(I1) aqueous thiocyanate complexes²⁹ except that the spin density transferred from Co(I1) to nitrogen is of an order of magnitude higher than that reported here for $Ln(III)$ SCN^- complexes.

The disagreement between the conclusions of this work and those of other authors^{$3-6$} concerning the nature of lanthanide(II1) thiocyanate complexes is possibly due to the difficulty to deduce this nature from the variation of thermodynamic constants. In the case of thiocyanato complexes, these constants are indeed very small $(\Delta H = 0.8 \text{ kcal mol}^{-1}, \Delta S = 4.6$ cal deg⁻¹ mol⁻¹⁶) and the experimental uncertainties are probably important as shown by the discrepancies between the values reported in different works.^{3,6} Moreover, the variations of thermodynamic constants result from a large number of effects so that the semiempirical rules³⁰ relating the sign of these constants to the nature of the complexes are not necessarily valid in the case of the weak thiocyanate complexes.

Registry No. $[Pr(SCN),(H,O),]^{+}$, 67905-00-4; $[Nd(SCN), (H_2O)_7$ ⁺, 67905-01-5; $[Eu(SCN)_2(H_2O)_7]$ ⁺, 67905-02-6; [Gd $(SCN)_2(H_2O)_7$ ⁺, 67905-03-7; [Tb(SCN)₂(H₂O)₇]⁺, 67905-04-8; $[Dy(SCN)₂(H₂O)₇]⁺$, 67905-05-9; $[Ho(SCN)₂(H₂O)₇]⁺$, 67905-06-0; $[Yb(SCN)_2(H_2O)_7]^+$, 67905-07-1; ¹³C, 14762-74-4; SCN⁻, 302-04-5.

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Comparison of the Electrochemical Properties and Electron Spin Resonance Spectra of Thio and Imino Chelates

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Received July **7,** *I978*

Electrochemical reduction and oxidation of square-planar nickel, palladium, and platinum complexes of the bidentate ligand $(NH)_2C_2(CN)^{2-}$ have been accomplished. The general electron-transfer series $MN_1^+ \rightleftharpoons MN_4^0 \rightleftharpoons NN_4^- \rightleftharpoons MN_4^2$ is detected, where N4 stands for **2** mol of the ligand. Cyclic voltammetry and phase-selective ac polarography establish the electron-transfer steps as rapid and reversible. Visible and near-IR spectra of the anions are greatly similar to those observed previously for the corresponding **MS4-** complexes, in which sulfur replaces NH in the chelate structure. Electron spin resonance spectra of the anions show that the half-filled orbital is highly delocalized over the chelate structure, and metal hyperfine splittings suggest that, in PtN₄⁻, the unpaired electron is more *ligand localized* than in PtS₄⁻. An unequivocal determination of the electronic ground states of these complexes could not be made, but the ESR spectra are more consistent with the **2A,** ground state than with the ${}^{2}B_{2g}$ ground state formed for the analogous dithiolate chelates.

Square-planar metal dithiolate complexes which undergo reversible electron-transfer reactions have been widely studied. Syntheses of analogous chelates with other donor groups (e.g., O or NH) have been reported,²⁻⁸ but characterization of the electron-transfer properties and electronic structures of these complexes has been less complete. Reductions of imino or oxo complexes are more difficult to achieve since these complexes

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have a lower electron affinity than the thio chelates^{$2-4$} and at

Table 1. Voltammetric Data for MN, Complexes

^{*a*} Volts vs. SCE. ^{*b*} Slope of plot of *-E* vs. log $[i/(i_d - i)]$ in mV. ^{*c*} Ratio of anodic to cathodic peak currents. ^{*d*} mV/s. *^{<i>e*} Separation</sub> between cathodic and anodic peak currents in mV.

least some appear subject to slower electron-transfer processes.⁹ The resultant difficulties in obtaining electron spin resonance (ESR) spectra of the paramagnetic monoanions derived from these species has hindered development of a quantitative description of the bonding in these chelates.

We now report electrochemical and spectroscopic studies of the reduction (and, to a lesser extent, oxidation) products of Ni, Pd, and Pt complexes formed utilizing the deprotonated form of the ligand diaminomaleonitrile, $(NH_2)_2C_2(CN)_2$. The complexes formed $(1, R = CN)$ are square planar around the metal^{5,10} and are congeners of the important maleonitriledithiolate (mnt) complexes $(2, R = \text{CN})$.^{11,12} The redox properties, electonic ground states, and degree of metal-ligand delocalization of the imino and thiolato complexes are then compared. Hereafter, these complexes are referred to as MN_4 and MS₄, respectively.

Experimental Section

Preparations. NiN₄ was prepared by a modification of the method of Miles et al.⁵ Use of anhydrous reagents allowed a substantial increase in yield (to about 30%). **A** nickel chloride-dimethoxyethane adduct was formed by refluxing anhydrous nickel chloride (2.5 mmol) in dry dimethoxyethane. After about 2 h, a yellow precipitate had formed. After the mixture was cooled to room temperature, 5.0 mmol of diaminomaleonitrile, dmn (Terramarine Bioresearch, recrystallized three times from ethyl acetate and vacuum dried), was added, and the solution was rendered basic by dropwise addition of 10.0 mmol of triethylamine. The intensely blue-green solution was filtered and evaporated, and the resulting powder was washed repeatedly with water. Final purification was accomplished using preparative-scale liquid chromatography employing a microstyragel column and tetrahydrofuran as the solvent. PdN_4^5 and PtN_4^{10} were prepared by the literature methods. Again, preparative liquid chromatography was used for final purification. $P dS_4^2$ and $P tS_4^2$ were prepared as their tetra-n-butylammonium salts by published procedures.¹³ The deuterated complex $\text{Ni}[(\text{ND})_2\text{C}_2(\text{CN})_2]_2$ was prepared by dissolving 20 mg of $\text{Ni} \text{N}_4$ - h_4 in 25 mL of acetone containing 4 mL of D_2O . After 24 h of stirring, the solution was filtered and evaporated. IR spectra of the product showed a strong N-D stretch (at ca. 2400 cm⁻¹) and only a low-intensity N-H stretch (ca. 3300 cm⁻¹). We estimate that about 90% deuteration was achieved. A ⁶¹Ni-enriched NiN₄ sample was prepared using 11.5 mg of ⁶¹Ni-enriched metal (90%, Oak Ridge). The nickel was dissolved in 1:1 nitric acid, heated to near dryness, and treated with a solution of 50 mg of dmn in *7* mL of methanol. After addition of a few drops of $Me₃N$, the solution turned deep green. After 1 h of stirring, the solution was evaporated and the residue washed several times with H_2O . The residue was dissolved in acetone, filtered, and evaporated to yield a small amount of 61 NiN₄.

Solvents and Supporting Electrolytes. Acetonitrile and N , N -dimethylformamide (DMF) were both spectroquality (MCB). The former was stirred for 24 h over $CaH₂$, degassed, and stored under vacuum. The latter was dried several days over Linde type 4A molecular sieves, decanted, and vacuum distilled from o-phthalic acid. 1,2-Dimethoxyethane (glyme) and tetrahydrofuran (THF) were dried

over either **LiAIH,** or Na, distilled, and stored under vacuum. Bu_4NPF_6 supporting electrolyte was prepared by mixing a slight excess of 65% HPF₆ with an alcoholic solution of Bu_4N (Southwestern Analytical Chemicals) and was recrystallized several times from 95% ethanol before vacuum drying. Electrolyte solutions were made 0.1 M in Bu_4NPF_6 . Mercury was chemically purified by the method of Booth and Jones¹⁴ before twice being distilled. All potentials are referred to the aqueous SCE.

Instrumentation. Princeton Applied Research Models 170 and **173** potentiostats were used for electrochemical experiments. Recording equipment, etc., was as previously described.¹⁵ Most electrochemical experiments were done using a nitrogen purge to eliminate oxygen, but bulk coulometry experiments were performed using the vacuum electrochemical cells described elsewhere.¹⁵ Electronic spectra of the ions were obtained by electrolytic reduction in a spectroelectrochemical cell which operated in the transmission mode. In these experiments a mercury pool cathode was aligned just below the optical light path in the spectrophotometer (Cary 14) and a solution was exhaustively electrolyzed at a potential appropriate for formation of the anion. ESR spectra was recorded with a Varian V-4502 spectrometer using a dual cavity with DPPH as a g-value standard. Paramagnetic MN_4 anions were prepared in one of two ways. In some cases, solutions of the anions produced by bulk electrolysis under nitrogen were syringed into ESR sample tubes. Typically, these samples were obtained after completion of coulometric experiments. Alternatively, solutions of $MN₄$ were electrolyzed in an intra muros cell¹⁶ directly in the microwave sample cavity and quick-frozen if desired. Computer simulation of ESR spectra and other computations utilized a Xerox Sigma-6 computer.

Results

Electrochemistry. The reduction of the MN_4 complexes at mercury electrodes was investigated by dc and ac polarography, cyclic voltammetry, and controlled-potential coulometry. Each complex is reduced in two one-electron steps to $MN_4^$ and MN_4^2 ⁻.

Formation of MN_4 **⁻ and** MN_4 **²⁻.** Figure 1 shows a dc polarogram of $NiN₄$ in dimethoxyethane. Two reduction waves of approximately equal height are accompanied by a small wave at about -0.7 V. This small wave is apparently due to adsorption of the complex and will be discussed later. The same pattern of waves was observed for all three complexes $[NiN₄, PdN₄, PtN₄ (Table I)].$ The two main waves were diffusion controlled, as shown by plots of limiting currents vs. square root of mercury column height which were linear and passed through the origin. Plots of dc potential vs. log $[i/(i_d$ $- i$)] gave slopes of -59 to -62 mV for both Faradaic waves of each complex. These values are consistent with each reduction being an electrochemically reversible one-electron transfer.

Phase-selective ac polarography confirmed both the electrochemical reversibility of the reductions and the adsorptive nature of the wave just negative of the first Faradaic process.

Figure 1. Dc polarogram of 3.0×10^{-4} M NiN₄ in 1,2-DME/ Bu_4NPF_6 (1,2-DME = 1,2-dimethoxymethane).

Figure 2 shows both the in-phase and out-of-phase components of an ac polarogram (at a frequency of 200 **Hz)** of NiN, in CH₃CN. The two symmetric peaks at -0.35 and -1.17 V correspond to the dc half-wave potentials of the two stepwise reductions to NiH_4^- and NiN_4^{2-} . From the observed currents, a standard rate constant for heterogeneous electron transfer, k_s , was calculated. The ratio of the in-phase to out-of-phase peak currents, cot ϕ , is related to k_s , D_o (the diffusion coefficient of the oxidized form $NiN₄$), and ω (2π times the ac frequency) by eq 1^{17} The diffusion coefficient for $NiN₄$

$$
\cot \phi = 1 + (2\omega D_0)^{1/2} / 2k_s \tag{1}
$$

was measured from our dc polarographic data as $D_0 = 1.20$ \times 10⁻⁵ cm²/s. Ac measurements at frequencies of 50, 100, and 200 Hz gave average values of $k_s = 0.25 \pm 0.03$ and 0.29 \pm 0.04 cm s⁻¹ for the reductions $\text{NiN}_4^0 \rightleftharpoons \text{NiN}_4^-$ and $\text{NiN}_4^ \Rightarrow$ NiN₄²⁻, respectively. These values correspond to highly reversible electron transfers giving no evidence of important structural changes during the charge transfer. Hence it is likely that the anions and dianions retain the square-planar structure. The electron-transfer rates are only slightly lower than those reported for oxidation and reduction of $\overline{MS_4}^{2-15,18,19}$ Thus, the reductions of these imino chelates are not subject to the slow electron transfers seen for some oxo chelates. 9,19

The ac polarogram of $NiN₄$ (Figure 2) also shows the asymmetric peak at ca. **-0.7** V, referred to previously, in which the out-of-phase (capacitive) component is considerably *greater* than the in-phase (resistive) component. This behavior is typical of adsorption-controlled processes^{20a} and has been observed for other metal complexes.20b The ac polarograms

Figure 3. Dc polarogram of 2.5 \times 10⁻⁴ M solution of PtN₄²⁻ made by electrolysis of PtN_4 at -1.6 V in 1,2-DME/Bu₄NPF₆.

of $PdN₄$ (in THF) were similar to those of the Ni complex, with Faradaic peaks at -0.43 and -1.08 V and an adsorption peak at -0.55 V. Adsorption problems were even more severe, however, for $PdN₄$, to the extent that the capacitive component to the ac wave was distorted in the region of the first reduction and k_s values could not be calculated for the PdN_4 reductions.

These complexes were also investigated by cyclic voltammetry (CV) at a hanging mercury drop electrode (HMDE). **A** variety of criteria established that the reductions were uncomplicated reversible charge transfers to give anions which were stable on the CV time scale (Table I). At slow scan rates, separations between the cathodic and anodic peaks were about 60 mV in agreement with value expected for a Nerstian system.21 Cathodic current functions were constant as a function of scan rate, also consistent with a charge transfer uncomplicated by coupled chemical reactions.

The long-term stability of the anions, required for synthetic or spectroscopic investigations, was investigated by controlled-potential coulometry. These studies established that all of the anions except NiN_4^{2-} are indefinitely stable under strictly anaerobic conditions. Typical of these results is the behavior of PtN_4^{2-} in dimethoxyethane. The polarogram after electrolysis of PtN_4 at -1.6 V shows (Figure 3) two waves resulting from oxidation of the *dianion* of PtN₄. The heights of the anodic waves were equal to the heights of the cathodic waves of the original $PtN₄$ solution, showing that complete conversion to PtN_4^2 had been effected.

The number of electrons transferred *(n)* and the anion stabilities were further studied by reversal coulometry. In this approach the compound is initially reduced for a given length of time, *t,* and then reoxidized for the same amount of time. The number of electrons transferred in the reduction, *n,* is calculated from a plot of log i vs. *t,* by reference to eq 2, in

$$
\log i = (-pt/2.30) + \log nFVC_0p \tag{2}
$$

which p is a constant (the mass-transfer coefficient), F is Faraday's constant, C_0 is the bulk concentration of the starting complex, and V is the solution volume. p is evaluated from the slope of the line, and n is calculated from the intercept.²² When the reduction product is stable and not subject to follow-up reactions, the numbers of coulombs consumed in the reduction (Q_i) and oxidation (Q_b) are theoretically related²² by eq 3. Representative values of *t* were 1000-2000 s.

$$
Q_{b}/Q_{f} = 1 - \exp(-pt) \tag{3}
$$

The values of *n* and Q_b/Q_f for $\text{NiN}_4^{0/-}$ (electrolysis potentials of -1.0 and 0 V) were 0.91 e and 0.95, respectively. The coulomb ratio compared very nicely with the value (0.96) measured by electronic integration. These data again show that the reduction of NiN_4 proceeds by one electron to a stable anion. Similar data were also obtained for the first PtN_4 reduction $[n = 1.01 \text{ e}, Q_b/Q_f = 0.38 \text{ (theory, eq 3)}, 0.38$ (measd)]. However, the same technique (potentials of -1.5)

 $a \text{ In cm}^{-1}$. $b \text{ log of absorbancy coefficient.}$

and -1.0 V) established that the NiN₄ dianion is unstable [n] = 1.20 e, $[Q_{b}/Q_{f}]$ (measd) is less than one-tenth of theory]. Exhaustive electrolysis of $PtN₄$ solutions at the potential of the second wave (-1.6 V) established that the dianion of $PtN₄$ is stable (vide ante, Figure 3), but no coulometry was performed on the second reduction. The Pd complex was investigated in less detail but it did appear that $PdN₄$ and PdN₄²⁻ were stable species, so *the only anion of the series which showed any sign of decomposition reactions was NiN42-.* Polarograms of solutions of $NiN₄$ after electrolysis at the second wave had very poorly defined waves. Thus, the decomposition of NiN_4^{2-} was not a simple oxidative regeneration of NiN_4 or NiN_4 . In contrast, air oxidation of solutions of all other anions resulted in quantitative recovery of the neutral MN_4 complex.

The oxidation of these complexes was briefly probed by voltammetry in acetonitrile at the rotating platinum electrode or by CV at a platinum bead. Oxidation waves were observed at $E_{1/2}$ potentials of +0.64 V (Ni), +0.43 V (Pd), and +0.98 V (\overrightarrow{Pt}). The NiN₄ oxidation was irreversible except at high scan rates $(>5 V/s)$. PtN₄⁺ was considerably more stable but follow-up reactions were still evident from slow CV scans. The Pd oxidation product appeared to film the electrode, and data on it were unreliable. Although these experiments were only of a preliminary nature, it does appear that cation radicals could be generated from these complexes at lower temperatures or in less reactive solvents (unfortunately, the complexes are virtually insoluble in CH_2Cl_2 , which is an excellent solvent for stabilizing oxidized compounds).

Optical Spectra of the Anions. Spectra of the neutral complexes and of the monoanions were recorded in the visible and near-infrared regions. Each neutral complex has an intense visible absorption band in the range 630-700 nm (Table 11). When solutions of the complexes were electrolyzed at the potential of the first reduction in a spectroelectrochemical cell, spectra of the monoanions were recorded. Each of the monoanions has an intense band in the near-IR region (e.g., Figure 4). Attempts to obtain clean spectra of the dianions were not successful. Electrolysis of $PtN₄$ at the potential of the second reduction gave spectra of the dianion contaminated by the monoanion. However, poorly resolved weaker bands (log $a_M \sim 3$) in the region 500–600 nm were attributed to Pt N_4^{2-} . Our attempts to obtain spectra of Pd N_4^{2-} (by either electrolysis or alkali-metal reductions) were frustrated by poor conversion of PdN_4^- to $PdN_4^2^-$.

The spectra of these complexes are very reminiscent of those of nickel-group dithiolenes. Absorptions of the neutral metal dithiolene complexes shift into the near-IR when the complexes are reduced to monoanions. These intense, low-energy transitions, which have been attributed to charge-transfer bands, are not seen in the spectra of the metal dithiolene

Figure 4. Visible-near-IR spectra of $NiN₄$ and $NiN₄$ recorded using an in situ electrolysis cell: broken line, before electrolysis; solid line, after electrolysis at -0.80 V.

Figure 5. ESR spectra of $NiN₄⁻$ in a glass: (a) $DMF/CH₃CN$ frozen-solution spectrum at 120 K; (b) computer-simulated spectrum using the parameters listed in Table I11 and line widths of *7* G for the low-field (g_1) feature and 6.3 G for the g_2 and g_3 components.

dianions. The similarity of the MN_4 and MS_4 spectra argues for qualitative similarities in the electronic structures of the two sets of complexes.

Electron Spin Resonance Spectra. As expected, the monoanions are paramagnetic. A solution containing electrogenerated $NiN₄⁻$ in a 2:1 DMF/CH₃CN solution gave a symmetric single-line ESR spectrum of 16-G peak-to-peak width. When this solution was frozen at 120 K, a spectrum of apparently axial symmetry resulted, with hyperfine splittings (hfs) visible on the g_{\perp} component (Figure 5a). Since ⁶¹Ni $(I = \frac{3}{2})$ has a very low natural abundance (1.2%), these splittings were ascribed to ^{14}N interactions arising from the ligands. Identical spectra were obtained in a 2-propanol glass (eliminating the possibility of $CH₃CN$ or DMF coordination as the origin of the splittings) and for a sample in which the hydrogen atoms of the ligand were exchanged with deuterium.

A. *g* **Values**

Table **111. ESR** Parameters for Nickel-Group dmn and mnt Complexes

a Data taken in CH,CN. Data taken from ref **23** and **26.** Data taken from ref *25.* Data taken from ref **24. e** Splittings in cm⁻¹. f Calculated; see text.</sup>

If the hfs were due to proton interactions, deuterium substitution would have led to a collapse of the hyperfine pattern since a ²H splitting is about one-sixth of the ¹H splitting $[g_N]$ $= 2.67$ (¹H), 0.41(²H)].

Extensive computer simulations of the g_{\perp} region were performed in order to assign the splittings (Table 111). The best fit (Figure 5b) was obtained with a slight deviation from axial symmetry in both the *g* value and the nitrogen hfs. Interaction with four equivalent nitrogens was calculated to be 5.8 \times 10⁻⁴ and 7.0 \times 10⁻⁴ cm⁻¹ in the g_2 and g_3 components, respectively. Although the simulation shows better resolution than the experimental spectrum on the high-field side of the region, this probably arises from line-width effects resulting from broadening of the nitrogen lines arising from higher *MI* values.

A frozen-solution spectrum of the Ni-61 enriched complex showed a well-resolved quartet along the g_{\parallel} region arising from the Ni hfs (Figure 6). There was no apparent change in the *g,* region, meaning that the Ni splitting along *g,* is very small. A similar pattern has been observed for $NiS₄$ ⁻.^{23,24}

The isotropic spectrum of $PdN₄$ was obtained from an acetonitrile solution of the anion transferred under nitrogen from the spectroelectrochemical cell to a quartz flat cell. Nitrogen hfs was clearly observed as nine lines due to interaction with four equivalent nitrogens ($a_N = 4.2 \times 10^{-4}$ cm⁻¹, Figure 7). No satellites due to Pd hfs $(^{105}Pd, 22%$ abundant) were observed.

In a separate experiment, PdN4- was frozen at 77 **K** in a 2:l DMF/THF solution and the glassy spectrum recorded (Figure 8). A rhombic **g** tensor is apparent, and, as with NiN₄⁻, the largest metal splitting (14.9 \times 10⁻⁴ cm⁻¹) is ob-

Figure **7.** Fluid-solution ESR spectrum of **WN,** in **CH3CN/Bu4NPF6** at **300 K.**

Figure **8.** Frozen-solution ESR spectrum of **PdN4-** in DMF/THF at **71** K.

served along the low-field component. This splitting was measured from the satellite lines around the low-field *g* component. Four of the six ¹⁰⁵Pd lines $(I = 5/2)$ are resolved, and their intensities were measured as 3.5% of the $I = 0$ line (theory = 4.7%). In the g_2 region, nitrogen hfs was seen (four equivalent N, nine lines, $a_N = 7.7 \times 10^{-4}$ cm⁻¹) but no metal splittings were observed. Along the high-field *(g,)* direction, a complex hyperfine pattern was observed which we failed to decipher. However, using the assumptions that a_{N_1} (low field) $= 0$ and that all the N hfs are of the same sign, we can calculate that a_{N_3} = 3.9 \times 10⁻⁴ cm⁻¹, using the measured values of a_{N_2} and a_{iso} and eq 4.

$$
a_{\rm iso} = \frac{1}{3}(a_1 + a_2 + a_3) \tag{4}
$$

 $PtN₄$ ⁻ was the only complex in which the metal hfs was observed for all three principal directions. Figure 9 shows the spectrum obtained from a frozen (77 K) DMF/THF solution

of the monoanion. All three central lines of the rhombic **g** tensor are flanked by a doublet arising from the 195Pt splitting $(I = \frac{1}{2}$, 33.8% abundant). These satellites are all about the correct intensity *(25.5%)* compared to the unsplit parent line. A fluid-solution spectrum in CH₃CN allowed observation of the isotropic Pt splitting, which was quite large: 156×10^{-4} cm-'. All of the Pt splittings have the same sign, as we can tell by comparing a_{Pt} (iso) with the average of the three anisotropic splittings. If all three have the same sign. an average of 155×10^{-4} cm⁻¹ is calculated, almost identical with the observed isotropic splitting. No well-resolved nitrogen hfs were observed for this complex, perhaps because the line widths were generally larger than in the other complexes. Splittings were observed along the middle *g* component which must be ascribed to nitrogen fine structure. These can be partially observed in Figure 9, but other spectra using long scan times and expanded magnetic field axes were used to make actual measurements. A nitrogen splitting of about 8×10^{-4} cm⁻¹ seems to account for the splittings, but because of the poor resolution this number should not be taken too seriously.

Electronic Ground State Assignments

In using the ESR data to evaluate the nature of the half-filled orbital of the anions, we need to assign the molecular directions of the three principal **g** tensor components in the

Since only the spectra of randomly oriented glasses of the ions could be produced, we assigned directions by analogy to previous single-crystal studies of MS_4^- and MS_{4}^- complexes. Each of these studies showed that the g_1, g_2 , and g_3 components (low to high field) were due to g_{yy} , g_{xx} , and g_{zz} , respectively,²³⁻²⁷ and we assumed these directions in our data analysis.

Various ground states were considered for these complexes, using the general approach of Maki et $al.^{23}$ In this approach, the experimental observables (g values and metal hyperfine splittings) are related to the bonding parameters *P* and *K* through the configurational excitation energies C_i for each ground state. For example, equations for the ${}^{2}B_{2g}$ ground state, composed of a mixture of the metal d_{xz} orbital and ligand p_z orbitals, are given as eq 5-10.

$$
g_{xx} = g_e + 2C_4 \tag{5}
$$

$$
g_{yy} = g_e - 6C_1 - 2C_3 \tag{6}
$$

$$
g_{zz} = g_e - 2C_2 \tag{7}
$$

$$
A_{xx} = P[2C_4 - K + \frac{2}{7} + \frac{3}{7}(C_1 + C_2 - C_3)] \tag{8}
$$

$$
A_{yy} = P[-6C_1 - 2C_3 - K - \frac{4}{7} - \frac{3}{7}(C_2 - C_4)] \tag{9}
$$

$$
A_{zz} = P[-2C_2 - K + \frac{2}{7} + \frac{3}{7}(C_3 - C_1 - C_4)] \tag{10}
$$

In these equations *P* and *K* are related to the anisotropic and isotropic parts, respectively, of the hyperfine tensor and are defined by eq *11* and *12. g,* is the free electron g value

$$
P = g_{\rm e}g_{\rm N}\beta_{\rm e}\beta_{\rm N}\langle r^{-3}\rangle \tag{11}
$$

$$
K = \left[-\frac{\gamma_3 \chi g_e g_N \beta_e \beta_N}{hc a_o^3} \right] / p \tag{12}
$$

 (2.0023) , g_N and β_N are the nuclear g value and nuclear magneton, respectively, and $\langle r^{-3} \rangle$ is the inverse cube electron-nucleus distance. χ is related to the spin polarization of s electrons by unpaired d electrons.^{28,29} $C_1 = \xi / (E_{d,z} - E_{d_{xx}})$

Figure 9. Frozen-solution ESR spectrum of PtN₄⁻ in DMF/THF at 17 K.

Table IV. Calculated Bonding Parameters for MN₄⁻ and MS₄⁻ Complexes for Various Possible Ground States

	NiN_{4} -	NIS_4^-	PtN_{a}	PtS_4^- ^d	
		${}^2B_{2g}$			
K^a_{pb} P/P ^c	0.28 29 0.26	0.21 24 0.21	-1.49 107 0.24	-0.1 270 0.6	
		2A_g			
K^+ P P/P_0	0.28 29 0.26	0.20 $25 -$ 0.22	-1.29 116 0.26	-0.40 175 0.39	

^{*a*} Defined by eq 12 in text. ^{*b*} Defined by eq 11 in text; units of 10^{-4} cm⁻¹. CValue of $P_0 = 110 \times 10^{-4}$ cm⁻¹ (Ni); 445 \times 10⁻⁴ cm⁻¹ (Pt). *d* Data taken from ref 24.

and the other constants involve excitation energies between the d_{xz} and d_{yz} (C₂), $d_{x^2-y^2}$ (C₃), and d_{xy} (C₄) orbitals.

In this treatment, *P* and *K* values are calculated for each possible ground state, based on the observed g and *A* values and a set of equations like (5) - (10) appropriate for that ground state [these equations are compiled in ref *23* (note that definition of the *x* and *y* directions are interchanged between our work and that of ref 23 and that the sign of C_4 in eq 15b of ref 23 has been shown²⁴ to be incorrect]. The sign and magnitude of *P* calculated for each ground state is then compared with the value (P_0) calculated for the free metal ion. The reasonableness of the *K* and P/P_0 values are then used to judge whether or not a particular ground state should be considered further as a possibility to describe the electronic structure of the complex. The ratio P/P_0 represents the lowering of the metal anisotropic parameter due to metalligand covalency and can be roughly taken as the degree of metal contribution in the half-filled orbital.

We also considered the possibility of an ${}^{2}A_{g}$ ground state configuration arising as a hybrid of the $d_{x^2-y^2}$ and d_{z^2} orbitals, which are allowed to mix in D_{2h} symmetry. This was not originally considered in the treatment of the metal dithiolenes but has since been discussed by Schlupp²⁴ and by McGarvey.³⁰ This orbital can be defined as $a(d_{z^2}) + b(d_{x^2-y^2})$ and would give rise to considerable *in-plane* anisotropy. Using the equations of Schlupp,²⁴ we calculated the values of $a^2 = 0.75$ and $b^2 =$ 0.25 (for NiN_4^-) and $a^2 = 0.66$ and $b^2 = 0.34$ (for PtN_4^-) for this orbital.

The results of calculations of *P* and *K* for the two most probable ground states (${}^{2}B_{2g}$ and ${}^{2}A_{g}$) are given in Table IV. All other possible ground states gave unreasonable values for P/P_0 and/or *K* (e.g., a negative value of P/P_0). The results show that for either choice of ground state for MN_4^- , the metal only contributes about one-quarter to the half-filled orbital. Thus, both NiN_4^- and PtN_4^- have the unpaired electron delocalized throughout the chelate structure, just as does **NiS4-.** It does appear that $PtN₄$ has considerably less metal character

to the unpaired electron than $P₄$. In this regard, the Pt hfs, which are considerably larger in the spectrum of PtN_4^- , are at first sight misleading, because their magnitude is due primarily to a large *isotropic* contribution. This is manifested in the large *negative* values of K calculated for $PtN₄$ ⁻ (Table IV). McGarvey has shown that *K* should be *positive* for third-row transition metals if the s character arises from spin polarization of the s electrons by the metal d electrons. A negative value such as we observe must be ascribed to *direct mixing* of the metal s orbital into the electronic ground state. Since this mixing must be rather large (3-4% contribution), the negative *K* value would seem to favor the ${}^{2}A_{g}$ ground state, which allows for direct mixing of the metal's and hybrid orbitals in D_{2h} symmetry.

However, the *ligand* (nitrogen) hfs hold the key to assignment of the ground state. McGarvey has pointed out that a measurable in-plane nitrogen hfs should accompany the ${}^{2}A_{\alpha}$ ground state,³⁰ whereas the ² B_{2g} ground state should give a large out-of-plane nitrogen splitting and small in-plane splittings. Indeed, Schmitt and Maki²⁶ used ³³S splittings to choose between the B_{2g} (d_{xz}) and A_{g} (d_z²) ground states for NiS₄. They noted that the ³³S hyperfine tensor $(A_{zz} = 14)$ $\times 10^{-4}$ cm⁻¹, $A_{xx} = A_{yy} = 4 \times 10^{-4}$ cm⁻¹) was only consistent with the former (d_{xz}) description because the *unique* splitting was along the *z (high-field)* direction.

Assuming, as noted above, that g_1 , g_2 , and g_3 correspond to g_{yy} , g_{xx} , and g_{zz} , in our MN_4^- complexes the unique nitrogen hfs appears to be along the low-field (y) direction, and a measurable A_{xx} for nitrogen is seen along the g_2 direction.

This can be best seen in the spectrum of the Ni complex, which has a narrow low-field g_1 resonance. We estimated that splittings of as little as *2* G would be clearly visible on this component. On the other hand, much larger splittings (6-7 G) are observed in the $x(g_2)$ and $z(g_3)$ directions. Similarly, a large nitrogen splitting is observed for the x direction (g_2) component) for both PdN_4 ⁻ and PtN_4 ⁻. By use of the arguments of Schmitt and Maki²⁶ and of McGarvey,³⁰ these observed ligand splittings cannot be reconciled with the ${}^{2}B_{2g}$ ground state. One further possibility which we considered was that the large a_N values along the g_{xx} direction arose from a misalignment of the principal directions of the ligand hyperfine tensor and the **g** tensor. This could still have allowed assignment of the ${}^{2}B_{2g}$ ground state. However, calculations reported in the Appendix show that coordinate misalignments cannot account for the magnitude of the apparent $a_N(x)$ splitting. We conclude that the ${}^{2}A_{g}$ ground state is the most likely description of the electronic structure of $NiN₄$ and $PtN₄$.

It must be emphasized that this assignment cannot be unequivocal without confirmation that the **g** tensor directions are as assumed in the glassy spectrum. However, the high air sensitivity of the anions may preclude examinations of their single-crystal spectra.

As a final point, calculations may be carried out using the observed nitrogen splittings to estimate the nitrogen spin density. The s- and p-orbital spin densities on nitrogen can be calculated through eq 13-16.

$$
A_{\parallel} = A_{\text{iso}} + 2B \tag{13}
$$

$$
A_{\perp} = A_{\text{iso}} - B \tag{14}
$$

$$
C_{\rm s}^2 = A_{\rm iso}/A_{\rm calcd} \tag{15}
$$

$$
C_{\rm p}^2 = B/B_{\rm calcd} \tag{16}
$$

Thus, the nitrogen spin densities C_s^2 and C_p^2 are obtained by comparison of the isotropic and anisotropic parts of the nitrogen hyperfine tensor $(A_{\text{iso}}$ and *B*, respectively) with the calculated versions of those quantities which assume complete localization of the unpaired electron in either the s or p orbital.

Table V

Values of A_{calod} and B_{calod} depend on the values of $|\chi_{2s}(0)|^2$ and $\langle r^{-3} \rangle_{2p}$ chosen for nitrogen. We used the values calculated by Morton et al.,³¹ in which $A_{\text{calcd}} = 552$ G and $B_{\text{calcd}} = 17.2$ G. Coupled with our experimental values of $A_{\text{iso}} = 4.3$ G and *B* $= 2.1$ G, these equations yield $C_s^2 = 0.008$ and $C_p^2 = 0.12$. In order to carry out these calculations, we assumed that the unresolved $A(y)$ splitting was zero and that all the nitrogen splittings were of the same sign. These results place half of the unpaired spin directly on the four coordinating nitrogens and confirm that the bulk of the unpaired spin is on the ligand.

Comparison of MN4 and MS4

In their overall properties, the imino and thio chelates are rather similar. Each set of complexes can exist (albeit sometimes transiently) in four of the five states of the general electron-transfer series

$$
MX_4^+ \rightleftharpoons MX_4 \rightleftharpoons MX_4^- \rightleftharpoons MX_4^{2-} \rightleftharpoons MX_4^{3-}
$$

The main difference in electrochemical behavior stems from the lower electron affinity of the nitrogen complexes. Standard reduction potentials are about 1.5 V more positive for the $MS₄$ complexes.³² This makes the $MN_4^{2-}/3$ - couple electrochemically inaccessible because of the limited cathodic window (to -2.8 V) of our electrolytes. Conversely, it is the most oxidized form, MS_4 ⁺, which remains undetected for these MS_4 complexes with strong electron-withdrawing (cyano) substituents. Each set of complexes gives rapid, reversible electron transfers in nonaqueous media. The electronic distributions (metal vs. ligand) are qualitatively similar and result in close analogies in electronic spectra and in potential separations between successive redox steps.³² It does appear that there is higher localization on the metal in $P₄$ -than in $P₄$ -. Finally, there is sufficient evidence in favor of the ${}^{2}A_{g}$ ground state for $MN_{4}^$ that it should be given strong consideration in discussions of these complexes.

Acknowledgment. The authors are grateful to the National Science Foundation for the support of this work and are indebted to **A.** H. Maki for providing access to data prior to publication and to Professor Maki, R. L. Schlupp, and R. D. Bereman for helpful discussions.

Appendix

Consideration of Possible Misalignment of Principal Directions of Nitrogen Hyperfine Tensor and g Tensor. This could result in an unexpectedly high value of $a_N(2)$ because the *g2* direction would not lie along a minimum in the nitrogen hfs. Such an effect has been observed in the ESR spectrum of a related complex, $Ni[Se_2C_2(CF_3)_2]_2^{-24}$ In this situation, there are a family of solutions to the ligand hyperfine tensor which reproduce the observed hfs. 33 The solutions involve the A_{xx} and A_{yy} hyperfine components and the angle of noncoincidence, θ , of the in-plane ligand hfs and **g** value tensors.

In principle, one can choose the correct solution from the set of possible solutions because each solution predicts small differences in the spacing of the ligand hyperfine lines, and these predictions can be mated with experimental observations. This approach yielded a value of $\theta = 32{\text -}36^{\circ}$ for the nickel-selenium complex mentioned above.24

Since our ligand hfs was not sufficiently resolved to allow observation of small deviations from the normally equal spacing of the hyperfine components, we sought to answer only one question with this technique: was there *any* reasonable value of θ which would allow a very small value of A_{xx} to be consistent with the observed nitrogen splitting along the g_{xx} direction? An answer of *no* would mean that the d_{xz} ground state could be eliminated from consideration. Table V gives a representative sampling of the values of A_{yy} , A_{xx} , and θ which reproduce the observed values of $A_1(N)$ and $A_2(N)$. These results showed that (for calculations up to $\theta = 40^{\circ}$) the observed $A_2(N)$ splitting of 5.8 G was the *minimum* value possible for $A_{xx}(\mathbf{N})$. This fact argues against the ²B_{2g} ground state for $NiN₄$, and, by implication, $PtN₄$, which also has a large splitting for $A_2(N)$.

Registry No. NiN₄, 34303-07-6; PdN₄, 34303-08-7; PtN₄, $34341 - 53 - 2$; NiN₄⁻, 67904-96-5; PdN₄⁻, 67904-97-6; PtN₄⁻, 67904 -98-7; $PtN₄²⁻, 67904$ -99-8.

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Electronic Spectra and Resonance Raman Spectra of Mixed-Valence Linear-Chain Complexes of Platinum with 1,2-Diaminoethane

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Receired July 13, 1978

The resonance Raman spectra of the mixed-valence complexes $[Pt(en)X_2][Pt(en)X_4]$ and $[Pt(en)_2][Pt(en)_2X_2](ClO_4)_4$, $en = 1,2$ -diaminoethane and $X = Cl$ or Br, have been recorded at room temperature and at ca. 80 K by use of an excitation line whose wavenumber falls within the contours of the intense, axially polarized, mixed-valence bands of these complexes in the visible region. The resonance Raman spectra are characterized by the appearance of an intense progression v_1v_1 ,
where v_1 is the symmetric X-Pt^{IV}-X stretching mode of the platinum(IV) moiety. This progressi 14, and 10 for the above four complexes, respectively. Weak subsidiary progressions $\nu_n + \nu_1 \nu_1$ also appear in the spectra, where u_n (the enabling mode) is another Raman-active mode. From the observed progressions, the spectroscopic constants ω_1 and x_{11} are calculated for each complex. The excitation profile of the ν_1 band of all four complexes reaches a maximum near the band maximum of the axially polarized mixed-valence band of each complex, although in the case of [Pt(en)- $Br_2[Pt(en)Br_4]$ another maximum in the excitation profile probably occurs near 13 000 cm⁻¹. The results suggest the use of the resonance Raman technique to detect and to resolve overlapped electronic absorption bands.

Introduction

Recent resonance Raman (RR) spectroscopic studies on mixed-valence complexes have focused on the linear-chain halogen-bridged variety, viz., $[Pt(\text{et}n)_4][Pt(\text{et}n)_4X_2]X_4 \cdot nH_2O$, etn = ethylamine, $n = 4$ (X = Cl or Br) and $n = 0$ (X = Br or I),¹⁻³ and [trans-M(NH₃)₂X₂][trans-M'(NH₃)₂X₄], where **M** or $M' = Pt$ or Pd and $X = CI$ or Br.^{4,5} The present paper develops this subject further with a detailed study of the mixed-valence complexes $[Pt(en)X_2][Pt(en)X_4]$ and $[Pt-$

 $(en)_2$ [Pt(en)₂X₂](ClO₄)₄, where en = 1,2-diaminoethane and $X = C1$ or Br. The object of the investigation is to establish the nature of the resonance Raman spectra obtained when these complexes are excited with radiation which falls within the contour of the lowest energy mixed-valence transition in each case and, thus, to provide important spectroscopic information on the complexes. The neutral ones are known to possess the linear-chain halogen-bridged structure shown in Figure 1,⁶ the Pt^{II}-Pt^{IV} distance increasing in the order chloride